TABLE 2. Values of the Interplanar Distances of Zirconium under a Pressure $P=350\,$ kbar, and the Interplanar Distances of Its Modifications

oar, and the Interpressing P = 350 kbar					Structure of the ω phase[4]			Structure of new phase		
					d	hhl	I	d	hkl	. I
d	I	d	nat				1	1		
2.794 2.568 2.523 2.460 1.897 1.784 1.615 1.459 1.395 1.395 1.260 1.228 1.127 1.0858 1.029 0.9783 0.9534	med med v.s. s. w. med. v.w. s. med. v.w. w. med. v.w. m	2.798 2.573 2.459 1.894 1.616 1.463 1.399 1.368 1.350 1.287 1.229 1.0842 1.0360 0.9783	010 002 	33 33 	3.103 — 2.533 2.521 — 1.957 — 1.786 — 1.555 1.455 — — 1.323 1.317 — 1.129 — 1.062 1.037 1.029 1.008 0.9784 0.9525 0.9517	231	6 — 45 100 — 13 — 20 — 11 50 — 43 3 — 14 — 14 1 7 7 5 15 5 13 1 15	-	222 - 123 -	1000

Note. We use the following notation: d are the interplanar distances; hkl are the Miller indices; I is the intensity; v.s. is very strong; s is strong; med. is medium; w. is weak; v.w. is very weak.

As the table shows the new high-pressure phase we obtained in zirconium after a shock wave has passed through it is not the ω phase of [4]. This is indicated by the absence on x-ray diffraction photographs of the ten comparatively strong lines of the ω phase, and also the lack of correspondence of experimental and calculated intensities. The indexing of seven lines of the new phase showed that all of them correspond with great precision to a lattice with the structure of a body-centered cube with $a = 3.568 \pm 0.005 \,\text{Å}$. The density of the new phase $\rho = 6.656 \,\text{g/cm}^3$.

A high-pressure phase was also identified in titanium, but only in samples subjected to a pressure $p=350~\rm kbar$. On the x-ray diffraction photographs of these samples three new, not very strong lines are present. The absence of lines of a new phase in titanium samples subjected to a pressure $P=500~\rm kbar$, and their low-intensity in zirconium samples for the same pressures are apparently explained by their instability at the high residual temperatures of shock compression. Just as in zirconium, these lines can be ascribed to the strongest reflections from (011), (002), and (022) of a cubic phase with $a=3.27~\rm \AA$. Unfortunately, these lines coincide with (110), (021), and (220) of the ti-

tanium ω phase. The absence of other strong reflections of the ω phase compels us to presume that we succeeded in isolating the body-centered cubic high-pressure phase in titanium, as in zirconium.

The last column of Table 1 shows the parameters of the new metastable phases of Zr and Ti; they were, however, obtained at much lower temperatures.

LITERATURE CITED

- P. W. Bridgman, Proc. Am. Acad. Arts Sci., 76, 71 (1948).
- 2. P. W. Bridgman, Proc. Am. Acad. Arts Sci., 81, 165 (1952).
- 3. P. W. Bridgman, Proc. Am. Acad. Arts Sci., 76, 55 (1948).
- 4. J. C. Jamieson, Science, 3 (3562), 72 (1963).
- A. Gayaraman, W. Klement, and G. G. Kennedy, Phys. Rev., <u>131</u>, 644 (1963).
- V. L. Al'tshuler, I. M. Barkalov, I. N. Dulin,
 V. N. Zubarev, T. N. Ignatovich, and P. A.
 Yampol'skii, Khimiya Vysokikh Energii, 2, 88
 (1968).
- 7. L. I. Mirkin, Handbook of X-ray Structural Analysis [in Russian], GIFML, Moscow (1961).